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Growth and holographic storage properties of In:Ce:Cu:LiNbO₃ crystal

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Abstract

A series of In:Ce:Cu:LiNbO₃ crystals with different concentration of In_2O_3 were grown by the Czochralski method. The infrared transmission spectra and the photo-damage resistant ability of the crystals were measured. The OH⁻ absorption peak of In(3 mol%):Ce:Cu:LiNbO₃ crystal shifts to ultraviolet. The photo-damage resistant ability of In(3 mol%):Ce:Cu:LiNbO₃ crystal is about two orders of magnitude higher than those of pure LiNbO₃ and Ce:Cu:LiNbO₃ crystals. The diffraction efficiency, response time of In:Ce:Cu:LiNbO₃ crystals were tested by two-wave coupling experiment. The response time of In(3 mol%):Ce:Cu:LiNbO₃ crystal only one fourth of that of Ce:Cu:LiNbO₃. The mechanism of OH⁻ absorption peak shifting and photo-damage resistant ability enhancement were discussed. © 2007 Elsevier B.V. All rights reserved.

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1. Introduction

Lithium niobate (LiNbO₃) single crystal has been widely used for its unique electro-optic, acousto-optic and nonlinear optical properties [1,2]. The most important application is holographic storage because of its excellent photo-refractive properties [3,4]. Holographic storage devices using LiNbO₃ crystals as the storage substrate can achieve large capability holographic storage. Nowadays, people adopt non-volatile holographic storage to maintain the information for a long time in this photo-refractive crystal. The two-photon fixed method is a new technique to achieve the non-volatile holographic storage. Two dopants with different energy levels are doped into LiNbO₃ crystal. The electrons can be excited from the deeper level to the shallower level by the high laser energy and all electrons in the shallow level (including excited ones from deep level) are used for storage. Using Mn:Fe:LiNbO3 and Ce:Cu:LiNbO3 crystals [5,6], the non-volatile holographic storage has been achieved by means of the two-photon fixed method.

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However, it is found that Ce:Cu:LiNbO₃ has two shortcomings: low photo-damage resistant ability and long response time, which severely limit its applications [7,8]. The holographic storage properties of the LiNbO₃ crystal can be greatly improved by doping with MgO [9,10], ZnO [11,12], or In₂O₃ [13,14], Sc₂O₃ [15,16]. In this paper, we propose to dope In₂O₃, CeO₂ and CuO into LiNbO₃ to grow In:Ce:Cu:LiNbO₃ crystals. Ce and Cu are doped to act as the deeper and shallower energy levels in the non-volatile holographic storage. In³⁺ is doped to improve the photo-damage resistant ability and shorten the response time.

2. Experimental

2.1. Preparation of crystal samples

In(3 mol%):Ce:Cu:LiNbO₃, In(2 mol%):Ce:Cu:LiNbO₃, In(1 mol%):Ce: Cu:LiNbO₃ crystals were grown by the Czochralski method. The intermediate frequency furnace was used as the heater. The starting materials were Li₂CO₃, Nb₂O₅, In₂O₃, CeO₂ and CuO with the purity of 99.99%. The Li/Nb ratio of the initial melts to grow the congruent LiNbO₃ is 0.946. The concentrations of Ce and Cu are 0.1 and 0.03 wt.%, respectively. The concentrations of In³⁺ are 1, 2 and 3 mol%, respectively. To prepare the polycrystalline, these materials were mixed thoroughly, calcinated at 850 °C for 4 h, and sintered at 1150 °C for 4 h. The technology conditions for crystal growth were as follow: pulling rate of 2 mm h⁻¹, rotation rate of 20 rpm and temperature gradient of

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Table 1The compositions of the samples

Crystal	1#	2#	3#	4#	LiNbO3
In ³⁺ in melt (mol%)	1.0	2.0	3.0	0	0
In ³⁺ in crystal (mol%)	0.703	1.296	1.852	0	0
$k \text{ of } \text{In}^{3+}$	0.703	0.648	0.617	-	_
Ce ⁴⁺ in melt (wt.%)	0.1	0.1	0.1	0.1	0
Cu ²⁺ in melt (wt.%)	0.03	0.03	0.03	0.03	0



Fig. 1. Infrared transmission spectra of the samples.

30-40 °C cm⁻¹. The diameters and heights of these crystals were about 30 and 50 mm, respectively. The as-grown crystals were poled in another furnace where the temperature was held at about 1200 °C for 4 h with an applied direct current of 5 mA cm⁻². Then the crystals were cut into samples of 15 mm × 10 mm × 2 mm. Then all the wafers were polished to optical grade. The compositions of the samples were listed in Table 1.

2.2. Properties measurements

The concentrations of In^{3+} in the In:Ce:Cu:LiNbO₃ crystals were measured by inductively coupled plasma atomic-emission spectrometry (ICP-AES). The distribution coefficient (*k*) of In^{3+} was calculated by comparing the In^{3+} concentration in the crystal and the initial In^{3+} concentration in the melt. The results are listed in Table 1.

The infrared transmission spectra at room temperature ranging from 3000 to 4000 cm^{-1} were obtained with a Fourier IR spectrometer. Infrared transmission spectra at the partial range from 3400 to 3600 cm^{-1} are shown in Fig. 1.

The facula-distortion observing method was applied to measure the photodamage resistant ability of the crystals. The experimental set-up is shown in Fig. 2. The Ar⁺ laser (λ = 488.0 nm) was employed as the light source, and the output intensity of the beam could be regulated. The beam transited the light shed and convex lens and then focused on the sample. The focula could be observed on the screen. The polarization direction of the beam was along the



Fig. 2. Set-up for measurement of photo-damage resistant ability.

Table 2Photo-damage resistant abilities (*R*) of the samples

Sample	$R (\mathrm{W} \mathrm{cm}^{-2})$		
1#	9.3×10^{2}		
2#	1.4×10^{3}		
3#	2.6×10^{4}		
4#	1.7×10^{2}		
LiNbO ₃	1.9×10^{2}		



Fig. 3. Sketch of two-wave coupling optical path. M_1 , M_2 , M_3 : mirrors; BS: beam splitter; D_1 , D_2 : detector; S_1 , S_2 , S_3 : diaphragm; PC: computer.

Table 3 The holographic storage properties of In:Ce:Cu:LiNbO₃

Samples	η_{\max} (%)	$\tau_{w}(s)$	τ_{e} (s)	τ (s)
4#	62	220	610	134
1#	56	160	580	97
2#	53	118	426	74
3#	48	60	295	38

c-axis of the crystal. When the power intensity was low, there was no photodamage and the facula was round. When the power intensity achieved a certain value (the threshold value), the facula began to elongate along the *c*-axis of the crystal. The photo-damage occurred and the power density of the output beam was recorded by the power detector. This value is defined as the photo-damage threshold (*R*), which is a parameter to evaluate the photo-damage resistance ability of the crystal. The higher the photo-damage resistant ability is, the larger signal-noise ratio the crystal has. The results are shown in Table 2.

The diffraction efficiency and the response time of the In:Ce:Cu:LiNbO₃ crystals were measured by the two-wave coupling experiments. The optical path was shown in Fig. 3. The He–Ne laser ($\lambda = 632.8$ nm) was used as the light source in the experiment. The polarizing of the output beam was in the incident plane. The beam was separated two by the beam splitter. One is signal light I_s and the other is reference light I_r . The intensity of $I_r = I_s = 150$ mW cm⁻² and their crossing angle is $10^{\circ}-30^{\circ}$. The diffraction efficiency η is defined as the ratio of the diffraction light intensity I_{sd} to the transmitted light intensity I_{st} for the signal light. The response time τ is defined as the time from the light begins radiating to the diffraction efficiency η is up to $(1 - \exp(-1))$ of its maximum (η_{max}). The writing time τ_w and the erasing time τ_e can be also gotten. The experiment results are shown in Table 3.

3. Results and discussion

The Infrared transmission spectra of the samples are shown in Fig. 1. For LiNbO₃ crystal with ABO₃ form structure, H⁺ ions can come into the lattices and combines with O^{2-} ions to form OH⁻. The shifting of OH⁻ absorption peak reflects the location and threshold concentration of the doped ions. From Fig. 1, it can be seen that the OH⁻ absorption peaks of 1[#] and 2[#] samples are located at about 3486 cm^{-1} , while that of 3[#] sample shifts to 3507 cm^{-1} . The OH⁻ absorption peaks of the crystals with low In³⁺ concentration are similar as that of pure congruent LiNbO₃ crystal, which is located about 3484 cm^{-1} [17,18].

From the viewpoint of Li-vacancy model, there are Li⁺ vacancies and antisite Nb $(Nb_{Li})^{4+}$ defects in congruent pure LiNbO₃ crystal and the OH⁻ absorption spectra reflect the change of the defects structure[19-21]. The OH⁻ vibration absorption peak of pure congruent LiNbO₃ crystal located at 3484 cm⁻¹ mainly reflects the vibration condition of OH- around Li+ vacancies. In the case of In:Ce:Cu:LiNbO₃ crystals, when In³⁺ concentration is below its threshold concentration, In³⁺ ions occupy antisite Nb $(Nb_{Li})^{4+}$ or Li⁺ vacancies, and exist in the form of $(In_{Li})^{2+}$. They are repellent to H⁺ ions. In this case, the OH⁻ still vibrate around Li⁺ vacancies, so the absorption peaks of 1[#] and 2[#] samples are similar as that of pure LiNbO₃ crystal. When In³⁺ concentration is above its threshold concentration, the additional In³⁺ ions begin to occupy Nb-site and exist in the form of $(In_{Nb})^{2-}$. Because $(In_{Nb})^{2-}$ have stronger attraction to H⁺ than that of Li⁺ vacancies. So H⁺ ions collect around $(In_{Nb})^{2-}$ and the absorption peak of $3^{\#}$ sample mainly reflects the vibration condition of OH⁻ around $(In_{Nb})^{2-}$.

The results of the photo-damage resistant ability of In:Ce:Cu: LiNbO₃ crystals are shown in Table 2. It can be seen that the photo-damage resistant ability of In(3 mol%):Ce:Cu:LiNbO3 crystal is about two orders of magnitude higher than those of pure LiNbO₃ and Ce:Cu:LiNbO₃ crystals. From the analysis above, we know that the In³⁺ ions begin to occupy Nb-site when In^{3+} concentration is above its threshold concentration. At the same time, the Ce⁴⁺ and Cu²⁺ ions occupy Nb-site, too. They exit in the form of $(In_{Nb})^{2-}$, $(Ce_{Nb})^{-}$ and $(Cu_{Nb})^{3-}$. Their abilities to capture electrons decrease and the photoconductivity of In(3 mol%):Ce:Cu:LiNbO3 crystal increases rapidly. According to the Scalar expression, $\delta \Delta n \approx B \kappa \alpha I / \sigma$ [22], where B is the generalized electro-optical coefficient, κ is the glass constant, α is the optical absorption coefficient, I is the intensity of the light and σ is the conductivity ($\sigma = \sigma_d + \sigma_p, \sigma_d$ is the dark conductivity and σ_p is the photoconductivity, $\sigma_d \ll \sigma_p$, so $\sigma \approx \sigma_p$), the photorefractive ability of the LiNbO3 crystal increases with the decrease of the photoconductivity. So with the increasing of the photoconductivity, the photo-damage resistant ability of In(3 mol%):Ce:Cu:LiNbO₃ improves obviously.

The results of holographic storage properties of In:Ce:Cu: LiNbO₃ are shown in Table 3. As holographic storage materials, the crystals should possess high diffraction efficiency to represent the recorded images clearly. The shorter response time can improve the efficiency. From the results shown in Table 3, the response time of In($3 \mod \%$):Ce:Cu: LiNbO₃ is about one fourth of that of Ce:Cu:LiNbO₃ crystal, though the diffraction

efficiency has a little decrease. The In(3 mol%):Ce:Cu:LiNbO₃ crystal is more suitable for holographic storage.

4. Conclusions

The Ce:Cu:LiNbO₃ crystal has two drawbacks: low photodamage resistant ability and long response time. By doping In_2O_3 in Ce:Cu:LiNbO₃ crystal to grow In:Ce:Cu:LiNbO₃ crystal, the photo-damage resistant ability increase two orders of magnitude higher and the response speed increases four times higher than those of Ce:Cu:LiNbO₃ crystal, when the In^{3+} concentration reach its threshold concentration. In(3 mol%):Ce:Cu:LiNbO₃ crystal is a more potential holographic storage material compared with Ce:Cu:LiNbO₃ crystal.

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